Search for the High-Rank Symmetries in Subatomic Physics - K. Mazurek - J. Dudek

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Participants

• France: IPHC, Strasbourg

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• Poland: IFJ-PAN, Kraków, UMCS, Lublin

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- Results (2008 2011):
- 17 articles, 18 conference talks

<u>Outline</u>

Introduction: Macroscopic-Microscopic Method

- Jacobi/Poincaré Shape Transition
- Shape-Dependent Congruence Energy
- Octupole Vibrations

Summary

Macroscopic-Microscopic Method

$$M(Z, N; def) = ZM_{H} + NM_{n} - 0.00001433Z^{2.39} + E_{LSD}(Z, N; def) + E_{micr}(Z, N; def)$$

Macroscopic Energy: Lublin - Strasbourg Drop, Finite Range Liquid Drop Model Microscopic Energy

$$E_{micr} = E_{pair} + E_{shell}$$

'Pairing' Energy

$$E_{pair} = E_{BCS} + \overline{E}_{pc}$$

or with Particle Number Projection method

Macroscopic-Microscopic Method

Nuclear surface parametrization:

$$\mathcal{R}(artheta, arphi) = \mathsf{R}_0 \, \mathsf{c}(\{lpha\}) \sum_{\lambda, \mu} [1 + lpha_{\lambda, \pm \mu} \, \mathsf{Y}_{\lambda, \pm \mu}(artheta, arphi)]$$



 $\{\beta,\gamma\} \rightarrow \{\alpha_{20},\alpha_{22}\}$

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Jacobi shape transition

Poincaré shape transition



The ¹⁴²Ba with spin $I = 0 - 110 \hbar$ A. Maj et al. Int.J.Mod.Phys.E 19 (2010). Investigation of the temperature influence into the Poincare shape transition is in progress.

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Giant Dipole Resonances





Eksperiment: A. Maj et al., Nucl. Phys. A 731, 319 (2004).

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The GDR strenght function versus gamma energy ⁶³Mn 4.5 ⁷⁹Ga ¹⁴²Ba 3.5 142Batot σ (arb units) 3 2.5 L=90 2 1.5 1 0.5 0 0 5 10 15 20 25 30 35 40 Energy [MeV] 57Ti ⁸⁵Se ¹⁴²Ba 142Batot σ (arb units) 6 5 4 3 L=100

Energy [MeV]

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0 5 10 15 20 25 30 35 40

Shape-Dependent Congruence Energy



New definition of the shape dependence

 $\mathsf{E}_{cong}(\alpha_{20}) = \mathsf{W}_0(\mathsf{Z},\mathsf{N}) \cdot \mathsf{F}_{neck}(\alpha_{20})$

 $F_{neck}(\alpha_{20}) = 1 + \frac{1}{2} \left\{ 1 + \tanh\left[(\alpha_{20} - \alpha_{20}^0)/a_{neck} \right] \right\}.$

 $W_0(Z,N) = -C_0 \exp(-W|I|/C_0) \text{ with } I \equiv (N-Z)/A$

Wigner E, Phys. Rev. 51 (1937) 106,947

 Myers W D, Swiatecki W J, Nucl. Phys. A 81 (1966) 1, Nucl. Phys. A 612 (1997) 249.
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Octupole Vibrations





$$\hat{\mathbf{H}} = \frac{\hbar^2}{2\mathbf{B}} \frac{\delta^2}{\delta \alpha_{30}^2} + \frac{1}{2} \mathbf{C} \alpha_{30}^2$$

where B is the mass parameter and C the stiffness of the potential.

$$\langle \psi_{n}|\hat{H}|\psi_{m}
angle = e_{n}\delta_{nm}$$

$$\hat{\mathbf{H}} = \frac{\hbar^2}{2\mathbf{B}} \frac{\delta^2}{\delta \alpha_{30}^2} + \mathbf{V}(\alpha_{30})$$

A. Bohr, B. Mottelson, Nuclear Structure vol. II

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Octupole Vibrations

The first 1^- state is usually described as 0 - point motion vibration. For ^{142}Ba the energy of this level is 1.326 MeV.				
	ħω [MeV]	Β [ħ ² /MeV]	e ₀ [MeV]	
	0.5	30	1.89	
	1.0	110	1.48	
	1.5	450	1.2	



The mean values of the dynamical octupole deformation of 0 - state excitation - $\sqrt{\langle \psi_0 | \alpha_{30}^2 | \psi_0 \rangle}$ 1 - state excitation - $\sqrt{\langle \psi_1 | \alpha_{30}^2 | \psi_1 \rangle}$





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Summary and Conclusions

• Prediction, for the first time, of the Poincare shape transition in wide range of nuclear masses and spins.

• The new formula for the shape - dependent congruence energy is proposed and tested.

• The octupole vibrations in fast rotating nuclei are under investigation.